Monografías de la Real Academia de Ciencias de Zaragoza 30, 35–41, (2006).

# Some particularities of the collinear point $L_3$ in the RTBP

E. Barrabés<sup>1</sup>, J. M. Mondelo<sup>2</sup> and M. Ollé<sup>3</sup>

<sup>1</sup> Dept. Informàtica i Matemàtica Aplicada Universitat de Girona Avd. Lluís Santaló s/n, 17071 Girona, Spain.

<sup>2</sup> Dept. Matemàtiques, Universitat Autònoma de Barcelona, 08193 Bellaterra, Barcelona, Spain.

<sup>3</sup> Dept. Matemàtica Aplicada I. Universitat Politècnica de Catalunya, Diagonal 647, 08028 Barcelona, Spain. \*

#### Abstract

We are interested in studying the motion in a (big) neighborhood of the collinear equilibrium point  $L_3$  of the RTBP. We consider both the planar and spatial cases. Actually different kinds of invariant objects appear: periodic orbits, invariant tori, the associated invariant manifolds, collision manifolds and homoclinic and heteroclinic phenomena among others. In this communication, we just present some particularities of  $L_3$  and its 1-dimensional manifolds to show the difficulties that we have to cope with in order to give a global description of the motion in a global neighborhood of  $_3$ .

**Key words and expressions:** Restricted three-body problem, equilibrium point, periodic orbits, invariant tori, manifolds.

### 1 Introduction

Our framework is the circular RTBP, whose well known equations of motion depend on the mass parameter  $\mu \in (0, 1/2]$ , and in a rotating reference system are

$$\begin{aligned} \ddot{x} - 2\dot{y} &= \Omega_x, \\ \ddot{y} + 2\dot{x} &= \Omega_y, \\ \ddot{z} &= \Omega_z, \end{aligned}$$

where

$$\Omega = \frac{1}{2}(x^2 + y^2) + \frac{1 - \mu}{r_1} + \frac{\mu}{r_2} + \frac{1}{2}\mu(1 - \mu)$$

<sup>\*</sup>e-mail: barrabes@ima.udg.es, jmm@mat.uab.es, merce.olle@upc.edu



Figure 1.— Horseshoe motion

and  $r_1 = \sqrt{(x-\mu)^2 + y^2 + z^2}$ ,  $r_2 = \sqrt{(x-\mu+1)^2 + y^2 + z^2}$ . Such equations have 5 equilibrium points  $L_1, \ldots, L_5$  ( $L_1, L_2$  and  $L_3$  are the collinear ones, while  $L_4$  and  $L_5$  are the equilateral ones), the Jacobi first integral given by

$$C = 2\Omega - \dot{x}^2 - \dot{y}^2 - \dot{z}^2.$$

and the well known symmetry

$$(x, y, z, x', y', z', t) \to (x, -y, z, -x', y', -t).$$
(1)

(see [13]).

In this communication we will deal with the collinear equilibrium points, with special emphasis on the  $L_3$  point. Actually,  $L_3$  is responsible for the horseshoe motion (see Figure 1). This kind of motion has been observed in some interesting situations, like the asteroid motion of Janus and Epimetheus as coorbitals of Saturn (see [11], [6], [7], [8] and [12]), and the motion of some more recent near Earth Asteriods (see [3] and [5]). Concerning horseshoe motion related to  $L_3$ , we mention [2], where a mechanism that explains horseshoe motion in the planar RTBP, for  $\mu > 0$  small, is given from the  $\mu = 0$ case, and also [1] where many families of horseshoe periodic orbits (both in the planar and spatial RTBP) are computed. On the other hand, we point out that the study of  $L_3$  is also interesting from an academic point of view since it requires the development of new methodology. Actually the neighborhood of  $L_1$ ,  $L_2$  has been (and is being) systematically studied by both semi-analytical and numerical techniques (see [4] and [10]). But some of this methodology does not apply easily to the neighborhood of  $L_3$  (see below).

We finally remark that our aim would be the description (as global as possible) of a neighborhood (as large as possible) of  $L_3$  including homoclinic and heteroclinic phenomena, for different values of the mass parameter  $\mu$ . Therefore we need to compute, in particular, periodic orbits (PO), invariant tori and the invariant manifolds of both PO



Figure 2.— Two families of Lyapunov PO.



Figure 3.— A sample torus.

and tori. Along this talk, we will describe some specific particularities of  $L_3$  and some problems that appear when studying and doing numerical explorations in a (maybe large) neighborhood of  $L_3$ .

#### **2** Local/global behaviour around $L_3$

As it is well known,  $\operatorname{Spec} Df(0) = \{\pm i\omega_1, \pm i\omega_2, \pm\lambda\}$ , so  $L_3$  is a center×center×saddle type equilibrium point. Thus, a first insight to the local dynamics around  $L_3$  gives two families of non-linear Lyapunov periodic orbits (PO) associated with the two centers (see Figure 2) and a 2-parametric (cantorian) family of 2-dimensional tori (see Figure 3). Of course, as far as the Jacobi constant decreases, the amplitudes of both the PO and tori increase. Therefore, in order to describe the global dynamics around  $L_3$ , it will be necessary to explore, in particular, the possible homoclinic connections to the existing invariant objects:  $L_3$ , the PO and the tori.

Therefore, first the homoclinic connections to  $L_3$ , when  $\mu$  increases, is explored. Then,



Figure 4.— Horsehsoe-shaped invariant manifold of  $L_3$ .

the case of the Lyapunov families of PO, when  $\mu$  is fixed and  $C_J$  decreases, will be studied. In this case, it will be necessary to take into account the bifurcations that appear in both, planar and vertical Lyapunov families. Finally, we will consider the families of invariant tori.

#### **3** Difficulties around $L_3$

On one hand, let us say that the invariant 1-dimensional manifolds of  $L_3$  may be regarded as the skeleton of the 2-dimensional manifolds of the Lyapunov PO close to  $L_3$ (when decreasing the Jacobi constant) and therefore of the 3-dimensional manifolds of the tori associated with the Lyapunov PO. On the other hand, the specific particularities of  $L_3$  give rise to some numerical difficulties when analysing the homoclinic connections. Let us mention some of these particularities:

- For small  $\mu$ , both the unstable and stable manifolds of  $L_3$ ,  $W^u(L_3)$  and  $W^s(L_3)$ , return to the neighborhood of  $L_3$ , see Figure 4.
- Let us take small values of  $\mu$  and let us concentrate on planar homoclinic orbits to  $L_3$ . We will distinguish between symmetric and non-symmetric homoclinic orbits. Due to symmetry (1), if an invariant manifold of  $L_3$  has an orthogonal crossing with the  $\{y = 0\}$  axis (i.e. for a suitable time t, we have y(t) = x'(t) = 0), then it becomes a symmetric homoclinic connection of  $L_3$ . This homoclinic orbit will have a horseshoe shape for  $\mu$  small enough. Actually, if one considers the homoclinic connections of  $L_3$  with only one crossing with the x axis (the orthogonal one), there exists a sequence of values of  $\mu_k \to 0$  for which this is the case (see [9] and [2]).
- The loops of the invariant manifolds that appear close to  $L_3$  are inherited by the invariant manifolds of planar Lyapunov orbits. That is, if we decrease the Jacobi constant, the Lyapunov PO has 2-dimensional invariant manifolds with also a horse-shoe shape and loops close to the Lyapunov orbit (see Figure 5). Therefore, when



Figure 5.— Left: horsehoe-shaped invariant manifold of a Lyapunov PO. Right: zoom.



Figure 6.— Left: invariant manifold of  $L_3$  that also surrounds  $L_2$  and the small primary. Right: invariant manifold of  $L_3$  that collides with the small primary.

analysing the symmetric homoclinic orbits of the Lyapunov PO by means of the intersection between the invariant manifold and the y = x' = 0 plane at a certain crossing, a careful analysis that takes into account the number of crossings must be done. Alternatively, a method that gives the homoclinic connections, regardless the number of crossings, should be designed.

For bigger values of μ, the neighborhood of the small primary and L<sub>1,2</sub> play a role (see Figure 6 left) and we can even have manifolds that go to collision with the small primary, so a regularization must be done (we have used Levi-Civita coordinates). We show in Figure 6 right a collision manifold. We refer the interested reader to [2] where a systematic study of the invariant manifolds of L<sub>3</sub>, when varying μ, is carried out.

#### 4 Further work

Once the behaviour of the invariant manifolds of  $L_3$  is studied, the next step consists of consider the families of Lyapunov orbits and tori. For each one of these objects we will take into account the following two situations:

- Symmetric homoclinic connections. As before, due to the symmetry (1), it will be enough to take into account the branches of one of the invariant manifolds, W<sup>u</sup> or W<sup>s</sup>, associated to the object and their intersections with the section Σ = {y = 0}. As we have mention before, to develop a numerical method to find W<sup>u,s</sup> ∩ Σ without considering the number of crossings with the section will be an important tool.
- Non-symmetric homoclinic connections. In this case, we will look for intersections between both W<sup>u</sup> and W<sup>s</sup> invariant manifolds. One way to do this could be to find, firstly, the intersections of each manifold (separately) with a suitable section Σ (for example, x = 0.5 μ), and then the intersections of the manifolds on Σ. Nevertheless, again in this case will be interesting to develop a numerical tool in order to find homoclinic connections as zeros of a suitable function.

#### Acknowledgments

E. Barrabés and J.M. Mondelo are partially supported by the MCyT/FEDER grants BFM2003-09504-C02-01 and MTM2006-05849/Consolider. J.M. Mondelo is also supported by the MCyT/FEDER grant MTM2005-02139. M. Ollé is partially supported by the MCyT/FEDER grant MTM2006-00478.

## References

- E. Barrabés and S. Mikkola. Families of periodic horseshoe orbits in the restricted three-body problem. Astron. Astrophys., 432(3):1115–1129, 2005.
- [2] E. Barrabés and M. Ollé. Invariant manifolds of  $L_3$  and horseshoe motion in the restricted three-body problem. *Nonlinearity*, 19:2065–2089, 2006.
- [3] R. Brasser, K.A. Innanen, M. Connors, D. Veillet, P. Wiegert, S. Mikkola, and P.W. Chodas. Transient co-orbital asteroids. *Icarus*, 171(1):102–109, September 2004.
- [4] E. Canalias and J. Masdemont. Homoclinic and heteroclinic transfer trajectories between Lyapunov orbits in the Sun-Earth and Earth-Moon systems. Discrete and Continuous Dynamical Systems, 14(2):261–279, 2006.
- [5] M. Connors, P. Chodas, S. Mikkola, P. Wiegert, C. Veillet, and K. Innanen. Discovery of an asteroid and quasi-satellite in an earth-like horseshoe orbit. *Meteoritics & Planetary Science*, 37:1435–1441, 2002.

- [6] Josep M. Cors and Glen R. Hall. Coorbital periodic orbits in the three body problem. SIAM J. Appl. Dyn. Syst., 2(2):219–237 (electronic), 2003.
- [7] Stanley F. Dermott and Carl D. Murray. The dynamics of tadpole and horseshoe orbits. i. theory. *Icarus*, (48):1–11, 1981.
- [8] Stanley F. Dermott and Carl D. Murray. The dynamics of tadpole and horseshoe orbits. ii. the coorbital satellites of saturn. *Icarus*, (48):12–22, 1981.
- [9] J. Font. The role of homoclinic and heteroclinic orbits in two-degrees of freedom Hamiltonian systems. Ph. D. Thesis, Barcelona University, 1999.
- [10] G. Gómez, W. S. Koon, M. W. Lo, J. E. Marsden, J. Masdemont, S.D. Ross. Connecting orbits and invariant manifolds in the spatial restricted three-body problem. *Nonlinearity*, 17:1571–1606, 2004.
- [11] J. Llibre, M. Ollé. The motion of Saturn coorbital satellites in the restricted three-body problem, Astron. Astrophys, 378: 1087–1099, 2001.
- [12] F. Spirig, J. Waldvogel. The three-body problem with two samll masses: a singularperturbation approach to the problem os Saturn's coorbiting satellites. *Stability of the solar* system and its minor and Artificial Bodies, ed. V. Szebehely (Reidel), pp. 53–64, 1985.
- [13] V. Szebehely. Theory of orbits. Academic Press, 1967.